

# On Estimating the Quantum $\ell_\alpha$ Distance

Yupan Liu  

Graduate School of Mathematics, Nagoya University, Japan

Qisheng Wang  

School of Informatics, University of Edinburgh, UK

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## Abstract

We study the computational complexity of estimating the quantum  $\ell_\alpha$  distance  $T_\alpha(\rho_0, \rho_1)$ , defined via the Schatten  $\alpha$ -norm  $\|A\|_\alpha := \text{tr}(|A|^\alpha)^{1/\alpha}$ , given  $\text{poly}(n)$ -size state-preparation circuits of  $n$ -qubit quantum states  $\rho_0$  and  $\rho_1$ . This quantity serves as a lower bound on the trace distance for  $\alpha > 1$ . For any constant  $\alpha > 1$ , we develop an efficient *rank-independent* quantum estimator for  $T_\alpha(\rho_0, \rho_1)$  with time complexity  $\text{poly}(n)$ , achieving an *exponential* speedup over the prior best results of  $\exp(n)$  due to Wang, Guan, Liu, Zhang, and Ying (*IEEE Trans. Inf. Theory* 2024). Our improvement leverages efficiently computable *uniform* polynomial approximations of *signed* positive power functions within quantum singular value transformation, thereby eliminating the dependence on the rank of the states.

Our quantum algorithm reveals a dichotomy in the computational complexity of the QUANTUM STATE DISTINGUISHABILITY PROBLEM WITH SCHATTEN  $\alpha$ -NORM ( $\text{QSD}_\alpha$ ), which involves deciding whether  $T_\alpha(\rho_0, \rho_1)$  is at least  $2/5$  or at most  $1/5$ . This dichotomy arises between the cases of constant  $\alpha > 1$  and  $\alpha = 1$ :

- For any  $1 + \Omega(1) \leq \alpha \leq O(1)$ ,  $\text{QSD}_\alpha$  is BQP-complete.
- For any  $1 \leq \alpha \leq 1 + \frac{1}{n}$ ,  $\text{QSD}_\alpha$  is QSZK-complete, implying that no efficient quantum estimator for  $T_\alpha(\rho_0, \rho_1)$  exists unless  $\text{BQP} = \text{QSZK}$ .

The hardness results follow from reductions based on new rank-dependent inequalities for the quantum  $\ell_\alpha$  distance with  $1 \leq \alpha \leq \infty$ , which are of independent interest.

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## 1 Introduction

Closeness testing of quantum states is a central topic in quantum property testing [58], which aims to develop (efficient) quantum testers for properties of quantum objects. This problem is also closely related to verifying the functionality of quantum devices, such as  $Q_0$  and  $Q_1$ , which are commonly designed to prepare the respective  $n$ -qubit (mixed) quantum states  $\rho_0$  and  $\rho_1$ . The goal of (tolerant) quantum state testing is to design efficient quantum



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algorithms that test whether  $\rho_0$  is  $2/5$ -far from or  $1/5$ -close to  $\rho_1$  with respect to a given closeness measure. Notably, this problem generalizes classical (tolerant) distribution testing (see [20] and [35, Chapter 11]) from a non-commutative perspective.

When the “source codes” of distribution- or state-preparation circuits are given, a surprising correspondence was established between such closeness testing problems – measured by the  $\ell_1$  norm distance [66, 37] or entropy difference [36] – and interactive proof systems that admit statistical zero-knowledge (SZK). This correspondence links closeness testing problems to both complexity theory and cryptography. A similar correspondence was later identified in the quantum world: closeness testing of quantum states with respect to the trace distance (given by Schatten 1-norm) [77, 78], denoted by QSD, or the von Neumann entropy difference [10] was shown to be QSZK-complete.<sup>1</sup>

In contrast, when the closeness measure follows an  $\ell_2$ -norm-like definition, such as the Hilbert-Schmidt distance or the quantum linear entropy, the corresponding closeness testing problems are in BQP using the SWAP test [17, 27]. Taken together, these results reveal a dichotomy in the complexity of closeness testing: when the measure is  $\ell_1$ -norm-like, the problems are QSZK-hard and their query or sample complexities have *polynomial* dependence on the dimension or rank of the states; whereas for  $\ell_2$ -norm-like measures, the problems are contained in BQP and their query or sample complexities are *rank-independent*.

What about the closeness testing problems with respect to generalizations that approximate the trace distance or the von Neumann entropy? The quantum  $\ell_\alpha$  distance, defined as  $T_\alpha(\rho_0, \rho_1) := \frac{1}{2} \text{tr}(|\rho_0 - \rho_1|^\alpha)^{1/\alpha}$  via the Schatten  $\alpha$ -norm, generalizes both the trace distance ( $\alpha = 1$ ) and the Hilbert-Schmidt distance ( $\alpha = 2$ ). Similarly, the quantum  $q$ -Tsallis entropy  $S_q(\rho)$  extends both von Neumann entropy ( $q = 1$ ) and quantum linear entropy ( $q = 2$ ).

Interestingly, prior results show a divergence in behavior for closeness measures looser than the  $\ell_2$  norm: The closeness testing problem with respect to  $T_\alpha(\rho_0, \rho_1)$ , denoted by  $\text{QSD}_\alpha$  (see Definition 17), is in BQP only for *even* integer  $\alpha \geq 2$  via the Shift test [27]; while for odd integers  $\alpha \geq 3$ , the query and sample complexities generally depend on the rank [71]. However, the techniques in [27] yield BQP algorithms for estimating  $S_q(\rho)$  for *all* integer  $q \geq 2$ . A recent work [53] further explored the closeness testing problem with respect to  $S_q(\rho_0) - S_q(\rho_1)$ , and extended the observed dichotomy from integers – where the transition occurs between  $q = 1$  and  $q \geq 2$  – to a continuous setting, showing a sharp distinction between  $q = 1$  and any constant  $q > 1$ . These results naturally lead to an intriguing question:

► **Problem 1.** What is the computational complexity of the closeness testing problem with respect to  $T_\alpha(\rho_0, \rho_1)$ ? Does a similar dichotomy hold between  $\alpha = 1$  and constants  $\alpha > 1$ , or does the complexity vary largely depending on whether  $\alpha$  is even or odd?

Why study  $\ell_\alpha$  problems for possibly non-integer  $\alpha > 1$ ? The trace distance ( $\alpha = 1$ ) is a fundamental closeness measure of quantum states, capturing the maximum success probability of quantum state discrimination [41, 40] and playing a key role in applications such as the security of quantum key distribution [11, 63]. For  $\alpha > 1$ , such as  $\alpha = 1.001$ , the quantum  $\ell_\alpha$  distance provides a natural lower bound on the trace distance, and addressing Problem 1 could make this bound efficiently computable. Moreover, insights from  $\ell_\alpha$  problems have previously contributed to progress on well-studied  $\ell_1$  problems, as seen in [49].

Beyond their connections to  $\ell_1$  problems,  $\ell_\alpha$  problems for  $\alpha > 1$  are of independent interest. In classical scenarios, they have applications in machine learning (e.g., [46]), as well as in streaming and sketching algorithms (e.g. [43]). In quantum scenarios, the Hilbert-

<sup>1</sup> The QSZK containment of the closeness testing problem with respect to the trace distance, denoted by  $\text{QSD}[a(n), b(n)]$ , holds only in the polarizing regime  $a(n)^2 - b(n) \geq 1/O(\log n)$ , as shown in [77, 78]. A recent work [51] slightly improved the parameter regime for this containment.

Schmidt distance ( $\alpha = 2$ ) is widely used in quantum information theory (e.g., [42, 61]), and more recently, has been leveraged in designing near-term (variational) quantum algorithms (e.g., [5, 28]). Consequently, positive answers to Problem 1 may offer new opportunities to refine, extend, or develop techniques relevant to these areas.

A classical counterpart to Problem 1 was investigated in [69] nearly a decade ago. The main takeaway aligns with [53]: For constant  $\alpha > 1$ , the sample complexity for distinguishing whether  $\text{TV}_\alpha(D_0, D_1)$  is at least  $2/5$  or at most  $1/5$  is *independent* of the dimension of the probability distributions  $D_0$  and  $D_1$ ,<sup>2</sup> fewer samples are needed as  $\alpha$  increases. Classically, these upper bounds can be achieved by drawing a polynomial number of samples and computing the  $\ell_\infty$  norm distance between the resulting empirical distributions. However, this approach does not directly extend to the quantum world: (1) quantum states  $\rho_0$  and  $\rho_1$  are generally not simultaneously diagonalizable; and (2) even when they are, estimating their eigenvalues associated with the unknown common eigenbasis remains challenging. Addressing these challenges is central to resolving Problem 1, which is the focus of our work.

## 1.1 Main results

We begin by stating our first main theorem, which provides a positive answer to Problem 1 when  $\alpha$  lies in the range  $1 + \Omega(1) \leq \alpha \leq O(1)$ :

► **Theorem 2** (Quantum estimator for  $T_\alpha$ , informal). *Given quantum query access to the state-preparation circuits of  $n$ -qubit quantum states  $\rho_0$  and  $\rho_1$ , for any constant  $\alpha > 1$ , there is a quantum algorithm for estimating  $T_\alpha(\rho_0, \rho_1)$  within additive error  $1/5$  and query complexity  $O(1)$ . If the state-preparation circuits have  $\text{poly}(n)$ -size descriptions, then the algorithm's time complexity is  $\text{poly}(n)$ . Thus, for any constant  $\alpha > 1$ ,  $\text{QSD}_\alpha$  is in BQP.*

More precisely, for a given additive error  $\epsilon$ , the explicit query complexity of Theorem 2 is  $O(1/\epsilon^{\alpha+1+\frac{1}{\alpha-1}})$  (see Theorem 14). In combination with the sampler [74, 75], estimating  $T_\alpha(\rho_0, \rho_1)$  can be done using  $\tilde{O}(1/\epsilon^{3\alpha+2+\frac{2}{\alpha-1}})$  samples of  $\rho_0$  and  $\rho_1$  (see Theorem 16). Both upper bounds can be expressed as  $\text{poly}(1/\epsilon)$  for the regime  $1 + \Omega(1) \leq \alpha \leq O(1)$ . In addition, if the state-preparation circuits of  $\rho_0$  and  $\rho_1$  have size  $L(n) = \text{poly}(n)$ , then Theorem 2 implies a quantum algorithm with time complexity  $\tilde{O}(L/\epsilon^{\alpha+1+\frac{1}{\alpha-1}})$ , or equivalently,  $\text{poly}(n, 1/\epsilon)$ .

Previous quantum algorithms for estimating the quantum  $\ell_\alpha$  distance for constant  $\alpha > 1$  have all relied on its powered variant, specifically the powered quantum  $\ell_\alpha$  distance:

$$\Lambda_\alpha(\rho_0, \rho_1) := \frac{1}{2} \text{tr}(|\rho_0 - \rho_1|^\alpha) = 2^{\alpha-1} \cdot T_\alpha(\rho_0, \rho_1)^\alpha.$$

Thus, for  $1 < \alpha \leq O(1)$ , the estimates of  $T_\alpha(\rho_0, \rho_1)$  and  $\Lambda_\alpha(\rho_0, \rho_1)$  are polynomially related.

When  $\alpha > 1$  is an even integer, estimating  $T_\alpha(\rho_0, \rho_1)$  follows from a folklore result via the Shift test [27], using  $O(1/\epsilon)$  queries or  $O(1/\epsilon^2)$  samples.<sup>3</sup> However, for odd integer  $\alpha > 1$ , no efficient quantum algorithm is known in general. Closeness testing of quantum states with respect to  $T_\alpha(\rho_0, \rho_1)$  for  $\alpha = 3$ , with query complexity  $O(1/\epsilon^{3/2})$ , has been noted only in [32]. For general non-integer constants  $\alpha > 1$ , the quantum query complexity of estimating  $T_\alpha(\rho_0, \rho_1)$  was studied in [71], with polynomial dependence on the maximum rank  $r$  of  $\rho_0$  and  $\rho_1$ . A technical comparison of our approach with this result is provided in Section 1.2.

<sup>2</sup> The closeness measure  $\text{TV}_\alpha(D_0, D_1)$  represents the classical  $\ell_\alpha$  distance based on the  $\ell_\alpha$  norm and generalizes the total variation distance, which is recovered at  $\alpha = 1$ .

<sup>3</sup> The sample complexity was noted in [62, Equations (83) and (84)].

By combining our efficient quantum estimator for  $T_\alpha(\rho_0, \rho_1)$  in the regime  $1 + \Omega(1) \leq \alpha \leq O(1)$  (Theorem 2) with our hardness results for  $\text{QSD}_\alpha$  (Theorem 3), we identify a sharp phase transition between the case of  $\alpha = 1$  and constant  $\alpha > 1$ , addressing Problem 1. For clarity, we summarize our main theorems and the quantitative bounds on quantum query and sample complexities, derived from both our results and prior work, in Table 1.

■ **Table 1** Computational, query, and sample complexities of  $\text{QSD}_\alpha$  for  $1 \leq \alpha \leq O(1)$ .

		$\alpha = 1$	$1 < \alpha \leq 1 + \frac{1}{n^{1+\delta}}$	$1 + \frac{1}{n^{1+\delta}} < \alpha \leq 1 + \frac{1}{n}$	$1 + \Omega(1) \leq \alpha \leq O(1)$
$\text{QSD}_\alpha$		QSZK-complete <sup>(*)</sup> [77, 78]	QSZK-complete <sup>(*)</sup> Theorem 3(2)	QSZK-complete <sup>(*)</sup> Theorem 3(2)	BQP-complete Theorems 2 and 3(1)
Query	Upper Bound	$\tilde{O}(r/\epsilon^2)$ [72]	$\tilde{O}\left(r^{3+\frac{2}{\alpha}}/\epsilon^{4\alpha+2}\right)$ [71]	$\tilde{O}\left(r^{3+\frac{2}{\alpha}}/\epsilon^{4\alpha+2}\right)$ [71]	$O\left(1/\epsilon^{\alpha+1+\frac{1}{\alpha-1}}\right)$ Theorem 14
	Lower Bound	$\tilde{\Omega}(r^{1/2})$ [18]	$\tilde{\Omega}(r^{1/2})$ Theorem 27(2)	$\Omega(r^{1/3})$ Theorem 27(1)	$\Omega(1/\epsilon)$ Theorem 22(1)
Sample	Upper Bound	$\tilde{O}(r^2/\epsilon^5)$ [72]	$\text{poly}(r, 1/\epsilon)$ Noted in [71, Footnote 2]	$\text{poly}(r, 1/\epsilon)$ Noted in [71, Footnote 2]	$\tilde{O}\left(1/\epsilon^{3\alpha+2+\frac{2}{\alpha-1}}\right)$ Theorem 16
	Lower Bound	$\Omega(r/\epsilon^2)$ [60]	$\Omega(r/\epsilon^2)$ Theorem 28	$\Omega(r/\epsilon^2)$ Theorem 28	$\Omega(1/\epsilon^2)$ Theorem 22(2)

(\*) For any  $\alpha(n) \in [1, 1 + \frac{1}{n}]$ , the promise problem  $\text{QSD}_\alpha[a, b]$  is contained in QSZK only under the polarizing regime  $a(n)^2 - b(n) \geq 1/O(\log n)$ , which can be slightly improved when  $\alpha = 1$  (see Footnote 1). However, establishing containment in a complexity class typically requires the natural regime  $a(n) - b(n) \geq 1/\text{poly}(n)$ , as in Theorem 2.

Finally, we present our second main theorem, which addresses the computational hardness of  $\text{QSD}_\alpha$ , as outlined in Theorem 3. In this context,  $\text{PUREQSD}_\alpha$  refers to a restricted variant of  $\text{QSD}_\alpha$  (see also Definition 17), where the states of interest are pure.

► **Theorem 3** (Computational hardness of  $\text{QSD}_\alpha$ ). *The promise problem  $\text{QSD}_\alpha$  captures the computational power of the respective complexity classes in the corresponding regimes of  $\alpha$ :*

- (1) **Easy regimes:** For any  $1 \leq \alpha \leq \infty$ ,  $\text{PUREQSD}_\alpha$  is BQP-hard. As a corollary,  $\text{QSD}_\alpha$  is BQP-complete for  $1 + \Omega(1) \leq \alpha \leq O(1)$ .
- (2) **Hard regimes:** For any  $1 \leq \alpha \leq 1 + \frac{1}{n}$ ,  $\text{QSD}_\alpha$  is QSZK-complete, where the QSZK containment of  $\text{QSD}_\alpha[a, b]$  only holds for the polarizing regime  $a(n)^2 - b(n) \geq 1/O(\log n)$ .

## 1.2 Proof techniques: BQP containment for $\alpha$ constantly above 1

At a high level, Quantum Singular Value Transformation (QSVT) [34] implies that the main challenge in designing a quantum algorithm based on a smooth function – e.g., Grover search [38] and the OR function, or the HHL algorithm [39] and the multiplicative inverse function (see [57] for more examples) – reduces to finding an efficiently computable polynomial approximation. Once such an approximation is obtained, the algorithm follows directly using techniques from [34], with its efficiency determined entirely by the polynomial's properties.

Now we focus on quantum algorithms for estimating the powered quantum  $\ell_\alpha$  distance. We begin by reviewing [71] and then provide a high-level overview of our approach.

The quantum query complexity of estimating the quantum  $\ell_\alpha$  distance for non-integer  $\alpha$  was first considered in [71, Theorem IV.1]. Their approach begins with the identity

$$2\Lambda_\alpha(\rho_0, \rho_1) = \|\rho_0 - \rho_1\|_\alpha^\alpha = \text{tr}\left(|\nu_-|^{\alpha/2} \Pi_{\nu_+} |\nu_-|^{\alpha/2}\right),$$

where  $\nu_{\pm} = \rho_0 \pm \rho_1$  and  $\Pi_{\nu_+}$  denotes the projector onto the support subspace of  $\nu_+$ . According to this identity, they aim to prepare a quantum state that is a block-encoding of (normalized)  $|\nu_-|^{\alpha/2} \Pi_{\nu_+} |\nu_-|^{\alpha/2}$ .<sup>4</sup> To this end, they first prepare a quantum state that is a block-encoding of  $\Pi_{\nu_+}$ , and then perform a unitary operator that is a block-encoding of  $|\nu_-|^{\alpha/2}$  on it.<sup>5</sup> Finally, the (unnormalized) powered quantum  $\ell_\alpha$  distance,  $\Lambda_\alpha(\rho_0, \rho_1)$ , can be obtained by estimating the trace of  $|\nu_-|^{\alpha/2} \Pi_{\nu_+} |\nu_-|^{\alpha/2}$  using quantum amplitude estimation [16]. After the error analysis, their approach was shown to have query complexity  $\tilde{O}(r^{3+1/\{\alpha/2\}}/\epsilon^{4+1/\{\alpha/2\}}) = \text{poly}(r, 1/\epsilon)$ .<sup>6</sup> The dependence on the rank is *inherent* in the approach of [71], as they have to prepare a rank-dependent quantum state that is a block-encoding of  $\Pi_{\nu_+}$ , making the rank parameters unavoidable in the error analysis.

To overcome this technical issue, we utilize an identity different from theirs:

$$2\Lambda_\alpha(\rho_0, \rho_1) = \|\rho_0 - \rho_1\|_\alpha^\alpha = \text{tr}(\rho_0 \cdot \text{sgn}(\nu_-) \cdot |\nu_-|^{\alpha-1}) - \text{tr}(\rho_1 \cdot \text{sgn}(\nu_-) \cdot |\nu_-|^{\alpha-1}).$$

The idea is to estimate the terms  $\text{tr}(\rho_j \cdot \text{sgn}(\nu_-) \cdot |\nu_-|^{\alpha-1})$  for  $j \in \{0, 1\}$  individually, and then combine them to obtain an estimate of  $\Lambda_\alpha(\rho_0, \rho_1)$ . Our algorithm is sketched as follows:

1. Find a good approximation polynomial for  $\text{sgn}(x) \cdot |x|^{\alpha-1}$ .
2. Implement a unitary block-encoding  $U$  of  $\text{sgn}(\nu_-) \cdot |\nu_-|^{\alpha-1}$  using Quantum Singular Value Transformation (QSVT) [34] and Linear Combinations of Unitaries (LCU) [23, 15], given the state-preparation circuits of  $\rho_0$  and  $\rho_1$ .
3. Perform the Hadamard test [3] on  $U$  and  $\rho_j$  with outcome  $b_j \in \{0, 1\}$  for each  $j \in \{0, 1\}$ .
4. Estimate  $\Lambda_\alpha(\rho_0, \rho_1)$  by computing the expected value of  $b_0 - b_1$ .

Our algorithm is actually inspired by the trace distance estimation in [72], which corresponds to the case of  $\alpha = 1$ . However, the approach in [72] still has a rank-dependent query complexity of  $\tilde{O}(r/\epsilon^2)$ , compared to the  $\tilde{O}(r^5/\epsilon^6)$  in [71].<sup>7</sup> Nevertheless, we discover an approach for estimating the quantum  $\ell_\alpha$  distance with a *rank-independent* complexity as long as  $\alpha$  is constantly greater than 1. Specifically, we use the best uniform approximation polynomial  $P_d(x)$  (of degree  $d$ ) for the function  $\text{sgn}(x) \cdot |x|^q$ , as given in [31, Theorem 8.1.1]:

$$\max_{x \in [-1, 1]} |P_d(x) - \text{sgn}(x) \cdot |x|^q| \rightarrow \frac{1}{d^q}, \text{ as } d \rightarrow \infty.$$

Our use of the best uniform approximation by polynomials is inspired by the recent work [53] on estimating the  $q$ -Tsallis entropy of quantum states for non-integer  $q$ , where they used the best uniform approximation polynomial for  $x^q$  in the non-negative range  $[0, 1]$  (given in [67]). The difference is that in our case, we have to further consider the sign of  $x$ , thereby requiring the polynomial approximation to behave well in the negative part. It turns out that the polynomial approximation given in [31] is suitable for our purpose. Having noticed this, we then use the now standard techniques (used in [48, 53]) such as Chebyshev truncations and the de La Vallée Poussin partial sum (cf. [65]) to construct efficiently computable asymptotically best approximation polynomials such that

$$\max_{x \in [-1, 1]} \left| P(x) - \frac{1}{2} \text{sgn}(x) \cdot |x|^q \right| \leq \epsilon, \quad \max_{x \in [-1, 1]} |P(x)| \leq 1, \quad \text{and} \quad \deg(P) = O\left(\frac{1}{\epsilon^{1/q}}\right).$$

<sup>4</sup> See Definition 10 for the formal definition of block-encoding.

<sup>5</sup> This is because of the evolution of subnormalized density operators [71, Lemma II.2].

<sup>6</sup> Here,  $\{x\} := x - \lfloor x \rfloor$  denotes the fractional part of  $x$ .

<sup>7</sup> Some readers may wonder if our approach applies to trace distance estimation ( $\alpha = 1$ ) to remove the rank dependence in quantum query and sample complexities in [72]. However, the answer is generally no, as the rank dependence is intrinsic to trace distance estimation due to the polynomial dependence of the rank in the quantum query and sample complexities lower bounds (see [52, Section 2.2.2] in the full version).

Using this efficiently computable polynomial (with  $q = \alpha - 1$ ) and with further analysis, we can then estimate the quantum  $\ell_\alpha$  distance to within additive error  $\epsilon$  with the desired query upper bound in Theorem 2. Moreover, using the (multi-)sampler [74, 75], a quantum query-to-sample simulation, we can also achieve the desired sample upper bound.

### 1.3 Proof techniques: QSZK completeness for $\alpha > 1$ near 1

To establish the BQP- and QSZK-hardness results in Theorem 3, we reduce the promise problems QSD and PUREQSD ( $\alpha = 1$ ) to the corresponding promise problems QSD $_\alpha$  and PUREQSD $_\alpha$  for appropriate ranges of  $\alpha$ . The key technique underlying these reductions is the following rank-dependent inequalities that generalize the case of  $\alpha = 2$  from [24, 25]:

► **Theorem 4** ( $T_\alpha$  vs.  $T$ , informal). *For any states  $\rho_0$  and  $\rho_1$  and  $\alpha \in [1, \infty]$ , it holds that:*

$$2^{1-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1) \leq T(\rho_0, \rho_1) \leq 2(\text{rank}(\rho_0)^{1-\alpha} + \text{rank}(\rho_1)^{1-\alpha})^{-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1). \quad (1)$$

For  $\alpha = \infty$ , the inequalities hold in the limit as  $\alpha \rightarrow \infty$ .

The proof of Theorem 4 follows from considering orthogonal positive semi-definite matrices  $\varsigma_0$  and  $\varsigma_1$  satisfying  $\rho_0 - \rho_1 = \varsigma_0 - \varsigma_1$ , and analyzing their properties carefully.

We then illuminate the hardness results in Theorem 3:

- For the easy regime, Equation (1) becomes an equality when both  $\rho_0$  and  $\rho_1$  are pure states. This equality implies the BQP-hardness of PUREQSD $_\alpha$ , along with the query and sample complexity lower bounds for all  $1 \leq \alpha \leq \infty$ , establishing Theorem 3(1).
- For the hard regime, Equation (1) is *sensitive* to  $\alpha$ . Particularly, for  $\alpha = 1 + \frac{1}{n}$ , if quantum states  $\rho_0$  and  $\rho_1$  are  $\tau$ -far, meaning  $T(\rho_0, \rho_1) \geq \tau$ , it follows only that  $T_{\alpha=1+\frac{1}{n}}(\rho_0, \rho_1) \geq \tau/2$ . However, when  $\alpha \leq 1 + \frac{1}{n^{1+\delta}}$  for any arbitrarily small constant  $\delta$ , the same trace distance condition ensures that  $T_\alpha(\rho_0, \rho_1) \geq \tau$  as  $n \rightarrow \infty$ , leading to the QSZK hardness result in Theorem 3(2) and distinct query complexity lower bounds in Table 1.

Lastly, we explain the QSZK containment in the hard regime. Simply combining Theorem 4 and the QSZK containment of QSD from [77, 78] does not work, as the resulting QSZK containment of QSD $_\alpha[a, b]$  holds only for  $a(n)^2/2 - b(n) \geq 1/O(\log n)$ , which is even weaker than the polarizing regime defined in Footnote 1. To address this, we establish a *partial* polarization lemma for  $T_\alpha$  (Lemma 25), which ensures that for quantum states  $\rho_0$  and  $\rho_1$  where  $T(\rho_0, \rho_1)$  is either at least  $a$  or at most  $b$ , we can construct new quantum states  $\tilde{\rho}_0$  and  $\tilde{\rho}_1$  such that  $T_\alpha(\tilde{\rho}_0, \tilde{\rho}_1)$  is either at least  $\frac{1}{2} - \frac{1}{2}e^{-k}$  or at most  $1/16$ , as long as the parameters  $a$  and  $b$  are in the polarizing regime. Theorem 3(2) follows by combining this partial polarization lemma for  $T_\alpha$  with the polarization lemma for  $T$  in [77].

### 1.4 Discussion and open problems

While the quantum  $\ell_\alpha$  distance  $T_\alpha(\cdot, \cdot)$  and its powered version  $\Lambda_\alpha(\cdot, \cdot)$  are almost computationally interchangeable for  $1 \leq \alpha \leq O(1)$ , their behavior differs greatly when  $\alpha = \infty$ :

- The quantity  $T_\infty(\rho_0, \rho_1)$  corresponds to the largest eigenvalue  $\lambda_{\max}$  of  $(\rho_0 - \rho_1)/2$ . The associated promise problem QSD $_\infty$  is BQP-hard and contains in QMA.<sup>8</sup> However, establishing a BQP containment appears challenging, as  $(\rho_0 - \rho_1)/2$  does not directly admit an efficiently computable basis – unlike its classical counterpart in [69], which does.

<sup>8</sup> The verification circuit in the QMA containment simply follows from phase estimation [45], where a (normalized) eigenvector corresponding to  $\lambda_{\max}$  serves as a witness state.

- The quantity  $\Lambda_\infty(\rho_0, \rho_1)$  takes values in  $\{0, 1/2, 1\}$  for any states  $\rho_0$  and  $\rho_1$ , and is nonzero if and only if the states are orthogonal, with at least one being pure. Thus, even the pure-state-restricted variant of the associated promise problem,  $\text{PUREPOWEREDQSD}_\infty[1, 0]$ , is  $\text{C=P}$ -hard (see [52, Appendix A] in the full version). Here,  $\text{C=P} = \text{coNQP}$  [2, 80], a subclass of  $\text{PP}$  that provides a precise variant of  $\text{BQP}$ , ensuring acceptance for all *yes* instances.

This fundamental difference between these quantities raises an intriguing question on  $\text{QSD}_\infty$ :

- (a) What is the computational complexity of the promise problem  $\text{QSD}_\infty$ , defined by  $T_\infty(\cdot, \cdot)$ ? Can we show that  $\text{QSD}_\infty$  is also in  $\text{BQP}$ , or is it inherently more difficult?

Another open problem concerns quantitative bounds for  $\text{QSD}_\alpha$ :

- (b) Can the query and sample bounds in Table 1 be improved, particularly for the regime  $1 + \Omega(1) \leq \alpha \leq O(1)$ ? Moreover, can tight bounds be established when the states have *small support*, analogous to the classical case in [69, Table 1]?

## 1.5 Related works

Schatten  $p$ -norm estimation  $\text{tr}(|A|^p)$  of  $O(\log n)$ -local Hermitian  $A$  on  $n$  qubits to within additive error  $2^{n-p}\epsilon\|A\|^p$  for  $\epsilon(n) \leq 1/\text{poly}(n)$  and real  $p(n) \leq \text{poly}(n)$  was shown to be  $\text{DQC1}$ -complete in [19]. Given a unitary block-encoding of a matrix  $A$ , in [56], they presented a quantum algorithm that estimates the Schatten  $p$ -norm  $(\text{tr}(|A|^p))^{1/p}$  to relative error  $\epsilon$  for integer  $p$ , where a condition number  $\kappa$  satisfying  $A \geq I/\kappa$  is required for the case of odd  $p$ .

The query complexity of  $N$ -dimensional quantum state certification (i.e., determine whether two quantum states are identical or  $\epsilon$ -far) with respect to trace distance was shown to be  $O(N/\epsilon)$  in [32]. The query complexity of trace distance estimation was shown to be  $\tilde{O}(r^5/\epsilon^6)$  in [71] and later improved to  $\tilde{O}(r/\epsilon^2)$  in [72], where  $r$  is the rank of the quantum states, confirming a conjecture in [25] that low-rank trace distance estimation is in  $\text{BQP}$ . Both Low-rank trace distance and fidelity estimations are known to be  $\text{BQP}$ -complete [1, 72]. Based on the approach of [72], space-bounded quantum state discrimination with respect to trace distance was shown to be  $\text{BQL}$ -complete in [48]. In addition to trace distance, fidelity is another important measure of the closeness between quantum states. The query complexity of fidelity estimation was shown to be  $\tilde{O}(r^{12.5}/\epsilon^{13.5})$  in [76] and later improved to  $\tilde{O}(r^{5.5}/\epsilon^{6.5})$  in [71] and to  $\tilde{O}(r^{2.5}/\epsilon^5)$  in [33]. Recently, the query complexity of pure-state trace distance and fidelity estimations was shown to be  $\Theta(1/\epsilon)$  in [70] and was recently extended in [29] to estimating fidelity of a mixed state to a pure state.

In addition to the query complexity, the sample complexity has also been studied in the literature. In [7], the sample complexity of  $N$ -dimensional quantum state certification was shown to be  $\Theta(N/\epsilon^2)$  with respect to trace distance and  $\Theta(N/\epsilon)$  with respect to fidelity. The sample complexity of trace distance estimation is known to be  $\tilde{O}(r^2/\epsilon^5)$  in [72] and that of fidelity estimation is known to be  $\tilde{O}(r^{5.5}/\epsilon^{12})$ , where  $r$  is the rank of quantum states. The sample complexity of pure-state squared fidelity estimation is known to be  $\Theta(1/\epsilon^2)$  via the SWAP test [17], where the matching lower bound was given in [4]. Recently, the sample complexity of pure-state trace distance and fidelity estimations was shown to be  $\Theta(1/\epsilon^2)$  in [73], which was achieved by using the sampler in [75].

## 2 Preliminaries

We assume a fundamental knowledge of quantum computation and quantum information theory. For an introduction, we refer the reader to the textbook [59].

We adopt the following notations throughout the paper: (1)  $[n] := \{1, 2, \dots, n\}$ ; (2)  $\tilde{O}(f)$  denotes  $O(f \text{ polylog}(f))$ , while  $\tilde{\Omega}(f)$  denotes  $\Omega(f / \text{polylog}(f))$ ; and (3)  $|\bar{0}\rangle$  represents  $|0\rangle^{\otimes a}$  for integer  $a > 1$ . In addition, the Schatten  $\alpha$ -norm of a matrix  $A$  is defined as  $\|A\|_\alpha := (\text{tr}(|A|^\alpha))^{1/\alpha}$ , where  $|A| := (A^\dagger A)^{1/2}$ . For simplicity, we use the notation  $\|A\|$  to denote the operator norm (equivalently, the Schatten  $\infty$ -norm) of a matrix  $A$ .

## 2.1 Closeness measures for quantum states

We start by defining the trace distance and providing useful properties of this distance:

► **Definition 5** (Trace distance). *Let  $\rho_0$  and  $\rho_1$  be two quantum states that are mixed in general. The trace distance between  $\rho_0$  and  $\rho_1$  is defined by*

$$T(\rho_0, \rho_1) := \frac{1}{2} \text{tr}(|\rho_0 - \rho_1|) = \frac{1}{2} \text{tr}\left(\left((\rho_0 - \rho_1)^\dagger (\rho_0 - \rho_1)\right)^{1/2}\right).$$

Notably, the trace distance is a distance metric (e.g., [79, Lemma 9.1.8]).

Next, we define the quantum  $\ell_\alpha$  distance and its *powered* version, which generalize the trace distance ( $\alpha = 1$ ) using the Schatten norm. Notably, the quantum  $\ell_\alpha$  distance coincides with the Hilbert-Schmidt distance when  $\alpha = 2$ :

► **Definition 6** (Quantum  $\ell_\alpha$  distance and its powered version). *Let  $\rho_0$  and  $\rho_1$  be two quantum states that are mixed in general. The quantum  $\ell_\alpha$  distance  $T_\alpha(\cdot, \cdot)$  and its powered version  $\Lambda_\alpha(\cdot, \cdot)$  between  $\rho_0$  and  $\rho_1$  are defined as follows:*

$$T_\alpha(\rho_0, \rho_1) := \frac{1}{2} \|\rho_0 - \rho_1\|_\alpha \quad \text{and} \quad \Lambda_\alpha(\rho_0, \rho_1) := \frac{1}{2} \|\rho_0 - \rho_1\|_\alpha^\alpha.$$

Here, the Schatten  $\alpha$ -norm of  $\rho_0 - \rho_1$  is given by  $\|\rho_0 - \rho_1\|_\alpha := (|\rho_0 - \rho_1|^\alpha)^{1/\alpha}$ .

By the monotonicity of the Schatten norm, e.g., [6, Equation (1.31)], it holds that:

$$\forall \alpha \geq 1, \quad 0 \leq \Lambda_\alpha(\rho_0, \rho_1) \leq T_\alpha(\rho_0, \rho_1) \leq T(\rho_0, \rho_1) \leq 1.$$

As a corollary of the Davis convexity theorem [26], the quantum  $\ell_\alpha$  distance  $T_\alpha(\cdot, \cdot)$  also serves as a distance metric, whereas its powered version  $\Lambda_\alpha(\cdot, \cdot)$  does not. Further inequalities for the trace distance and the quantum  $\ell_\alpha$  distance are given in the full version, particularly in [52, Section 2.1].

Lastly, we require the following relationship for additive error estimation between the quantum  $\ell_\alpha$  distance and its powered version. The proof is provided in the full version, specifically in [52, Proposition 2.5].

► **Proposition 7** ( $T_\alpha$  vs. powered  $T_\alpha$ ). *The quantum  $\ell_\alpha$  distance  $T_\alpha(\cdot, \cdot)$  and its powered version  $\Lambda_\alpha(\cdot, \cdot)$  are related through the equality  $T_\alpha(\rho_0, \rho_1) = 2^{\frac{1}{\alpha}-1} \cdot \Lambda_\alpha(\rho_0, \rho_1)^{\frac{1}{\alpha}}$ . Accordingly, if  $x$  is an estimate of  $\Lambda_\alpha(\rho_0, \rho_1)$  to within additive error  $\epsilon$ , then  $2^{\frac{1}{\alpha}-1} \cdot x^{\frac{1}{\alpha}}$  serves as an estimate of  $T_\alpha(\rho_0, \rho_1)$  to within additive error  $2^{\frac{1}{\alpha}-1} \cdot \epsilon^{\frac{1}{\alpha}}$ .*

## 2.2 Closeness testing of quantum states via state-preparation circuits

We begin by defining the closeness testing of quantum states with respect to the trace distance, denoted as QSD $[a, b]$ ,<sup>9</sup> and two variants of this problem:

<sup>9</sup> While Definition 8 aligns with the classical counterpart of QSD defined in [66, Section 2.2], it is slightly less general than the definition in [77, Section 3.3]. Specifically, Definition 8 assumes that the input length  $m$  and the output length  $n$  are *polynomially equivalent*, whereas [77, Section 3.3] allows for cases where the output length (e.g., a single qubit) is *much smaller* than the input length.

► **Definition 8** (Quantum State Distinguishability Problem, QSD, adapted from [77, Section 3.3]). Let  $Q_0$  and  $Q_1$  be quantum circuits acting on  $m$  qubits (“input length”) and having  $n$  specified output qubits (“output length”), where  $m(n)$  is a polynomial function of  $n$ . Let  $\rho_i$  denote the quantum state obtained by running  $Q_i$  on state  $|0\rangle^{\otimes m}$  and tracing out the non-output qubits. Let  $a(n)$  and  $b(n)$  be efficiently computable functions. Decide whether:

- Yes: A pair of quantum circuits  $(Q_0, Q_1)$  such that  $T(\rho_0, \rho_1) \geq a(n)$ ;
- No: A pair of quantum circuits  $(Q_0, Q_1)$  such that  $T(\rho_0, \rho_1) \leq b(n)$ .

Furthermore, we denote the restricted version, where  $\rho_0$  and  $\rho_1$  are pure states, as PUREQSD.

In this work, we consider the *purified quantum access input model*, as defined in [77], in both white-box and black-box scenarios:

- **White-box input model:** The input of the problem QSD consists of descriptions of polynomial-size quantum circuits  $Q_0$  and  $Q_1$ . Specifically, for  $b \in \{0, 1\}$ , the description of  $Q_b$  includes a sequence of polynomially many 1- and 2-qubit gates.
- **Black-box input model:** In this model, instead of providing the descriptions of the quantum circuits  $Q_0$  and  $Q_1$ , only query access to  $Q_b$  is allowed, denoted as  $O_b$  for  $b \in \{0, 1\}$ . For convenience, we also allow query access to  $Q_b^\dagger$  and controlled- $Q_b$ , denoted by  $O_b^\dagger$  and controlled- $O_b$ , respectively.

In addition to query complexity, defined within the black-box input model, *sample complexity* refers to the number of copies of quantum states  $\rho_0$  and  $\rho_1$  needed to accomplish a specific closeness testing task. Useful lemmas on computational hardness and quantitative lower bounds for query and sample complexities are available in the full version, specifically in [52, Section 2.2].

## 2.3 Polynomial approximations

We now present useful tools on best uniform polynomial approximations. In addition, definitions and lemmas on Chebyshev expansion and truncations can be found in the full version, particularly in [52, Section 2.3.2].

Let  $f(x)$  be a continuous function defined on the interval  $[-1, 1]$  that we aim to approximate using a polynomial of degree at most  $d$ . We define  $P_d^*$  as a *best uniform approximation* on  $[-1, 1]$  to  $f$  of degree  $d$  if, for any degree- $d$  polynomial approximation  $P_d$  of  $f$ , it holds that  $\max_{x \in [-1, 1]} |f(x) - P_d^*(x)| \leq \max_{x \in [-1, 1]} |f(x) - P_d(x)|$ .

The best uniform (polynomial) approximation of positive (constant) powers  $|x|^\alpha$  was first established by Serge Bernstein [14, 13]. However, the focus here is on the best uniform approximation of *signed* positive powers  $\text{sgn}(x)|x|^\alpha$ , as stated in Lemma 9. This result is often attributed to Bernstein’s work (see, e.g., [68, Equation (10.2)]), and a proof of a more general version can be found in [31, Theorem 8.1.1].

► **Lemma 9** (Best uniform approximation of signed positive powers, adapted from [31, Theorem 8.1.1]). For any positive real (constantly large) order  $\alpha$ , let  $P_d^* \in \mathbb{R}[x]$  be the best uniform polynomial approximation for  $f(x) = \text{sgn}(x)|x|^\alpha$  of degree  $d = \left\lceil (\beta_\alpha/\epsilon)^{1/\alpha} \right\rceil$ , where  $\beta_\alpha$  is a constant depending on  $\alpha$ . Then, for sufficiently small  $\epsilon$ ,  $\max_{x \in [-1, 1]} |P_d^*(x) - f(x)| \leq \epsilon$ .

## 2.4 Quantum algorithmic toolkit

In this subsection, we recap the quantum singular value transformation. Additional quantum algorithmic tools, including useful quantum algorithmic subroutines and the quantum sampler can be found in the full version, specifically in [52, Sections 2.4.2 and 2.4.3].

### 2.4.1 Quantum singular value transformation

We start by defining block-encoding:

► **Definition 10** (Block-encoding). *A linear operator  $A$  on an  $(n+a)$ -qubit Hilbert space is said to be an  $(\alpha, a, \epsilon)$ -block-encoding of an  $n$ -qubit linear operator  $B$ , if  $\|\alpha(|0\rangle^{\otimes a} \otimes I_n)A(|0\rangle^{\otimes a} \otimes I_n) - B\| \leq \epsilon$ , where  $I_n$  is the  $n$ -qubit identity operator and  $\|\cdot\|$  is the operator norm.*

Then, we state the quantum singular value transformation:

► **Lemma 11** (Quantum singular value transformation, [34, Theorem 31]). *Suppose that unitary operator  $U$  is a  $(\alpha, a, \epsilon)$ -block-encoding of Hermitian operator  $A$ , and  $P \in \mathbb{R}[x]$  is a polynomial of degree  $d$  with  $|P(x)| \leq \frac{1}{2}$  for  $x \in [-1, 1]$ . Then, we can implement a quantum circuit  $\tilde{U}$  that is a  $(1, a + 2, 4d\sqrt{\epsilon/\alpha} + \delta)$ -block-encoding of  $P(A/\alpha)$ , by using  $O(d)$  queries to  $U$  and  $O((a+1)d)$  one- and two-qubit quantum gates. Moreover, the classical description of  $\tilde{U}$  can be computed in deterministic time  $\text{poly}(d, \log(1/\delta))$ .*

## 3 Efficient quantum algorithms for estimating quantum $\ell_\alpha$ distance

In this section, we present efficient quantum algorithms for estimating the quantum  $\ell_\alpha$  distance  $T_\alpha(\rho_0, \rho_1)$  when  $\alpha \geq 1 + \Omega(1)$ . These algorithms utilize either queries to state-preparation circuits or samples of the states  $\rho_0$  and  $\rho_1$ . The core of our approach is an *efficient* uniform approximation to *signed* positive constant power functions (Lemma 12), which provides a uniform error bound over the entire interval  $[-1, 1]$ .

This uniform polynomial approximation enables a query-efficient quantum algorithm for estimating  $T_\alpha(\rho_0, \rho_1)$  through its powered version  $\Lambda_\alpha(\rho_0, \rho_1)$ , as shown in Theorem 14. As a result, we establish a BQP containment of the promise problem  $\text{QSD}_\alpha$  defined in Section 4. Additionally, by leveraging the multi-sampler in [73], we devise a sample-efficient quantum algorithm for estimating  $T_\alpha(\rho_0, \rho_1)$ , detailed in Theorem 16.

### 3.1 Efficient uniform approximations of signed positive powers

Using the averaged Chebyshev truncation specified in [52, Section 2.3.2] in the full version, we provide an *efficiently computable* uniform polynomial approximation of *signed* positive constant powers (see also [52, Lemma 3.1]):

► **Lemma 12** (Efficient uniform polynomial approximation of signed positive powers). *Let  $\alpha$  be a positive real (constantly large) number. For any  $\epsilon \in (0, 1/2)$ , there is a degree- $d$  polynomial  $P_d \in \mathbb{R}[x]$ , where  $d = \left\lceil (\beta'_\alpha/\epsilon)^{1/\alpha} \right\rceil$  and  $\beta'_\alpha$  is a constant depending on  $\alpha$ , that can be deterministically computed in  $\tilde{O}(d)$  time. For sufficiently small  $\epsilon$ , it holds that:*

$$\max_{x \in [-1, 1]} \left| \frac{1}{2} \text{sgn}(x) |x|^\alpha - P_d(x) \right| \leq \epsilon \quad \text{and} \quad \max_{x \in [-1, 1]} |P_d(x)| \leq 1.$$

More specifically, the degree- $d$  uniform approximation of signed positive powers (see Lemma 9) ensures that the degree- $(2d - 1)$  averaged Chebyshev truncation (given in [52, Equation 2.3]) achieves almost the same approximation error bound, with efficiently computable Chebyshev coefficients. The complete proof is available in the full version, specifically in [52, Section 3.1].

## 3.2 Quantum $\ell_\alpha$ distance estimation for constantly large $\alpha > 1$

### 3.2.1 Query-efficient quantum algorithm for estimating powered $T_\alpha$

We now present efficient quantum query algorithms for estimating  $\Lambda_\alpha(\rho_0, \rho_1)$  and  $T_\alpha(\rho_0, \rho_1)$ , as presented in [52, Lemma 3.2] in the full version, and the complete proof is provided in [52, Section 3.2.1]:

► **Lemma 13** (Powered quantum  $\ell_\alpha$  distance estimation via queries). *Suppose that  $Q_0$  and  $Q_1$  are unitary operators that prepare purifications of mixed quantum states  $\rho_0$  and  $\rho_1$ , respectively. For every constantly large  $\alpha \geq 1 + \Omega(1)$ , there is a quantum query algorithm that estimates  $\Lambda_\alpha(\rho_0, \rho_1)$  to within additive error  $\epsilon$  by using  $O(1/\epsilon^{1+\frac{1}{\alpha-1}})$  queries to  $Q_0$  and  $Q_1$ .*

By combining Proposition 7 with Lemma 13 for additive error  $\epsilon^\alpha$ , we obtain a quantum query algorithm for estimating  $T_\alpha(\rho_0, \rho_1)$  when  $\alpha \geq 1 + \Omega(1)$  is constantly large, as also stated in [52, Theorem 3.3] in the full version:

► **Theorem 14** (Quantum  $\ell_\alpha$  distance estimation via queries). *Suppose that  $Q_0$  and  $Q_1$  are unitary operators that prepare purifications of mixed quantum states  $\rho_0$  and  $\rho_1$ , respectively. For every constantly large  $\alpha \geq 1 + \Omega(1)$ , there is a quantum query algorithm that estimates  $T_\alpha(\rho_0, \rho_1)$  to within additive error  $\epsilon$  by using  $O(1/\epsilon^{\alpha+1+\frac{1}{\alpha-1}})$  queries to  $Q_0$  and  $Q_1$ .*

### 3.2.2 Sample-efficient quantum algorithm for estimating powered $T_\alpha$

We proceed by describing efficient quantum sample algorithms for  $\Lambda_\alpha(\rho_0, \rho_1)$  and  $T_\alpha(\rho_0, \rho_1)$ , as stated in Lemma 15 and in [52, Lemma 3.4] in the full version. Our sample algorithms are obtained by combining the quantum query algorithm from Lemma 13 with the sampler [74, 75]. The corresponding explanatory framework is presented in [52, Algorithm 1], and the detailed proof is given in [52, Section 3.2.2].

► **Lemma 15** (Powered quantum  $\ell_\alpha$  distance estimation via samples). *For every constantly large  $\alpha \geq 1 + \Omega(1)$ ,  $\Lambda_\alpha(\rho_0, \rho_1)$  can be estimated to within additive error  $\epsilon$  on a quantum computer by using  $\tilde{O}(1/\epsilon^{3+\frac{2}{\alpha-1}})$  samples of  $\rho_0$  and  $\rho_1$ .*

By combining Proposition 7 with Lemma 15 for additive error  $\epsilon^\alpha$ , we obtain a quantum sample algorithm for estimating  $T_\alpha(\rho_0, \rho_1)$  when  $\alpha \geq 1 + \Omega(1)$  is constantly large, as also stated in [52, Theorem 3.5] in the full version:

► **Theorem 16** (Quantum  $\ell_\alpha$  distance estimation via samples). *For every constantly large  $\alpha \geq 1 + \Omega(1)$ , there is a quantum sample algorithm that estimates the quantum  $\ell_\alpha$  distance  $T_\alpha(\rho_0, \rho_1)$  to within additive error  $\epsilon$  by using  $\tilde{O}(1/\epsilon^{3\alpha+2+\frac{2}{\alpha-1}})$  samples of  $\rho_0$  and  $\rho_1$ .*

## 4 Hardness and lower bounds for $\alpha$ constantly above 1

We begin by introducing a generalization of QSD from [77], where the trace distance is replaced with the quantum  $\ell_\alpha$  distance as the closeness measure:

► **Definition 17** (Quantum State Distinguishability Problem with Schatten  $\alpha$ -norm,  $\text{QSD}_\alpha$ ). *Let  $Q_0$  and  $Q_1$  be quantum circuits acting on  $m$  qubits (“input length”) and having  $n$  specified output qubits (“output length”), where  $m(n)$  is a polynomial function of  $n$ . Let  $\rho_i$  denote the quantum state obtained by running  $Q_i$  on state  $|0\rangle^{\otimes m}$  and tracing out the non-output qubits. Let  $a(n)$  and  $b(n)$  be efficiently computable functions. Decide whether:*

- Yes: A pair of quantum circuits  $(Q_0, Q_1)$  such that  $T_\alpha(\rho_0, \rho_1) \geq a(n)$ ;
- No: A pair of quantum circuits  $(Q_0, Q_1)$  such that  $T_\alpha(\rho_0, \rho_1) \leq b(n)$ .

Moreover, we denoted the restricted version, where  $\rho_0$  and  $\rho_1$  are pure states, as  $\text{PUREQSD}_\alpha$ .

In the remainder of this section, we establish rank-dependent inequalities between the quantum  $\ell_\alpha$  distance and the trace distance (Theorem 18) in Section 4.1. These inequalities facilitate reductions that demonstrate the BQP hardness (Theorem 21) and derive quantitative lower bounds on queries and samples (Theorem 22) for PUREQSD $_\alpha$  in Section 4.2.

#### 4.1 Rank-dependent inequalities between $T_\alpha$ and the trace distance

We generalize the rank-dependent inequalities between the (squared) Hilbert-Schmidt distance and the trace distance, as demonstrated in [24, Appendix G] and [25, Theorem 1] for the case of  $\alpha = 2$ , to all  $1 \leq \alpha \leq \infty$ :

► **Theorem 18** ( $T_\alpha$  vs.  $T$ ). *Let  $\rho_0$  and  $\rho_1$  be quantum states. The following holds:*

(1) *For any  $\alpha$  in the range  $1 \leq \alpha < \infty$ ,*

$$2^{1-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1) \leq T(\rho_0, \rho_1) \leq 2(\text{rank}(\rho_0)^{1-\alpha} + \text{rank}(\rho_1)^{1-\alpha})^{-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1).$$

(2) *For  $\alpha = \infty$ ,  $2 \cdot T_\infty(\rho_0, \rho_1) \leq T(\rho_0, \rho_1) \leq 2 \min\{\text{rank}(\rho_0), \text{rank}(\rho_1)\} \cdot T_\infty(\rho_0, \rho_1)$ .*

It is worth noting that Item 1 and Item 2 in Theorem 18 are consistent, specifically

$$\lim_{\alpha \rightarrow \infty} (\text{rank}(\rho_0)^{1-\alpha} + \text{rank}(\rho_1)^{1-\alpha})^{-\frac{1}{\alpha}} = \min\{\text{rank}(\rho_0), \text{rank}(\rho_1)\}.$$

Additionally, the inequalities in Theorem 18 sharpen the inequalities between the trace norm and the Schatten norm (see, e.g., [6, Equation (1.31)]):

$$\forall 1 \leq p \leq \infty, \quad \|A\|_p \leq \|A\|_1 \leq r_A^{1-1/p} \cdot \|A\|_p. \quad (2)$$

By considering the maximum rank of  $\rho_0$  and  $\rho_1$ , we can derive a simplified form of Theorem 18 for convenience:

► **Corollary 19** ( $T_\alpha$  vs.  $T$ , simplified). *For any quantum states  $\rho_0$  and  $\rho_1$ , the following holds:*

$$\forall 1 \leq \alpha < \infty, \quad 2^{1-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1) \leq T(\rho_0, \rho_1) \leq (2 \max\{\text{rank}(\rho_0), \text{rank}(\rho_1)\})^{1-\frac{1}{\alpha}} \cdot T_\alpha(\rho_0, \rho_1).$$

Moreover, for pure quantum states, Theorem 18 yields the following equality:

► **Corollary 20** ( $T_\alpha = T$  for pure states). *For any pure states  $|\psi_0\rangle\langle\psi_0|$  and  $|\psi_1\rangle\langle\psi_1|$ , we have:*

$$\forall 1 \leq \alpha \leq \infty, \quad 2^{1-\frac{1}{\alpha}} \cdot T_\alpha(|\psi_0\rangle\langle\psi_0|, |\psi_1\rangle\langle\psi_1|) = T(|\psi_0\rangle\langle\psi_0|, |\psi_1\rangle\langle\psi_1|).$$

The proof of Theorem 18 (see also [52, Theorem 4.2] in the full version) relies on the positive semi-definite matrices  $\varsigma_0$  and  $\varsigma_1$ , defined for each  $b \in \{0, 1\}$  as  $\varsigma_b := \frac{1}{2}((-1)^b(\rho_0 - \rho_1) + |\rho_0 - \rho_1|)$ . These matrices are orthogonal and satisfy the relation  $\rho_0 - \rho_1 = \varsigma_0 - \varsigma_1$ . The complete proof is provided in [52, Section 4.1].

#### 4.2 Computational hardness and lower bounds

We first state the computational hardness result of PUREQSD $_\alpha$  with  $1 \leq \alpha \leq \infty$ , obtained by a reduction from PUREQSD, which is BQP hard [64], as stated in Theorem 21 and in [52, Theorem 4.5]. The proof can be found in the full version, specifically in [52, Section 4.2].

► **Theorem 21** (PUREQSD $_\alpha$  is BQP-hard). *For any  $1 \leq \alpha \leq \infty$  and  $n \geq 2$ , it holds that:*

$$\text{PUREQSD}_\alpha \left[ 2^{\frac{1}{\alpha}-1} \cdot (1 - 2^{-n}), 2^{\frac{1}{\alpha}-1-n} \right] \text{ is BQP-hard.}$$

Next, we provide lower bounds on the query and sample complexities for  $\text{PUREQSD}_\alpha$  by reducing to those for  $\text{PUREQSD}$  in [70], as presented in Theorem 22 and in [52, Theorem 4.6]. The proof is provided in the full version, particularly in [52, Section 4.2].

► **Theorem 22** (Quantitative lower bounds for  $\text{PUREQSD}_\alpha$ ). *For any  $1 \leq \alpha \leq \infty$  and  $0 < \epsilon < 2^{\frac{1}{\alpha}-2}$ , there exist  $n$ -qubit pure states  $|\psi_0\rangle$  and  $|\psi_1\rangle$  such that deciding whether  $T_\alpha(|\psi_0\rangle\langle\psi_0|, |\psi_1\rangle\langle\psi_1|)$  is at least  $\epsilon$  or exactly 0 requires:*

(1) **Queries:** *In the purified quantum access model, the quantum query complexity is  $\Omega(1/\epsilon)$ .*

(2) **Samples:** *The quantum sample complexity is  $\Omega(1/\epsilon^2)$ .*

## 5 Quantum $\ell_\alpha$ distance estimation for $\alpha > 1$ near 1

In this section, we establish that  $\text{QSD}_\alpha$  is QSZK-complete for  $1 \leq \alpha \leq 1 + \frac{1}{n}$  (see also [52, Theorem 5.1] in the full version), extending the prior result that  $\text{QSD}$  ( $\alpha = 1$ ) is QSZK-complete, as shown in [77]:

► **Theorem 23** ( $\text{QSD}_\alpha$  is QSZK-complete for  $\alpha > 1$  near 1). *Let  $a(n)$  and  $b(n)$  be efficiently computable functions such that  $0 \leq b < a \leq 1$ . Then, for any  $1 \leq \alpha \leq 1 + \frac{1}{n}$ , it holds that:*

*For any  $a(n)^2 - b(n) \geq 1/O(\log n)$ ,  $\text{QSD}_\alpha[a, b]$  is in QSZK.*

*Moreover,  $\text{QSD}_\alpha[a, b]$  is QSZK-hard if  $a(n) \leq 1/2 - 2^{-n^\tau - 1}$  and  $b(n) \geq 2^{-n^\tau - \frac{1}{n+1}}$  for every constant  $\tau \in (0, 1/2)$  and sufficiently large integer  $n$ .*

The main challenge in proving Theorem 23 is to establish a QSZK containment of  $\text{QSD}_\alpha$  under the polarizing regime  $a(n)^2 - b(n) \geq 1/O(\log n)$ .<sup>10</sup> A direct approach, combining the inequalities between  $T$  and  $T_\alpha$  (Corollary 19) with the QSZK containment of  $\text{QSD}$  from [77, 78], only yields a QSZK containment of  $\text{QSD}_\alpha[a, b]$  under a *weaker* regime,  $a(n)^2/2 - b(n) \geq 1/O(\log n)$ . To circumvent this, we provide a (partial) polarization lemma for  $T_\alpha$  (Lemma 25), which enables us to achieve the desired QSZK containment in Theorem 23.

The remainder of this section establishes the QSZK containment of  $\text{QSD}_\alpha$  in Section 5.1. We then show the QSZK hardness of  $\text{QSD}_\alpha$  (Theorem 26) and derive quantitative lower bounds on query complexity (Theorem 27) and sample complexity (Theorem 28) in Section 5.2.

### 5.1 QSZK containment via a partial polarization lemma for $T_\alpha$

► **Theorem 24** ( $\text{QSD}_\alpha$  is in QSZK). *Let  $a(n)$  and  $b(n)$  be efficiently computable functions satisfying  $0 \leq b < a \leq 1$ . Then, the following holds:*

*For any  $\alpha \in \left[1, 1 + \frac{1}{n}\right]$  and any  $a(n)^2 - b(n) \geq \frac{1}{O(\log n)}$ ,  $\text{QSD}_\alpha[a, b]$  is in QSZK.*

The proof of Theorem 24 (see also [52, Theorem 5.2] in the full version) can be found in [52, Section 5.1]. The main technical tool is a *partial* polarization lemma for  $T_\alpha$  (see also [52, Lemma 5.3]), which ensures that any pair  $(a, b)$  within the polarizing regime can be made *constantly* separated:

<sup>10</sup>Notably, similar to the classical cases [12], by reducing to the QUANTUM JENSEN-SHANNON DIVERGENCE PROBLEM (QJSP) or the MEASURED QUANTUM TRIANGULAR DISCRIMINATION PROBLEM (MEASQTDP) introduced in [51], the QSZK containment of  $\text{QSD}_\alpha$  holds slightly beyond the polarizing regime.

► **Lemma 25** (A partial polarization lemma for  $T_\alpha$ ). *Let  $Q_0$  and  $Q_1$  be quantum circuits that prepare quantum states  $\rho_0$  and  $\rho_1$ , respectively. There is a deterministic procedure that, given an input  $(Q_0, Q_1, a, b, k)$  where  $a(n)^2 - b(n) \geq 1/O(\log n)$ , outputs new quantum circuits  $\tilde{Q}_0$  and  $\tilde{Q}_1$  that prepare the states  $\tilde{\rho}_0$  and  $\tilde{\rho}_1$ , respectively. The resulting states satisfy that:*

$$\begin{aligned} \text{For any } \alpha \in [1, 1 + 1/n], \quad T_\alpha(\rho_0, \rho_1) \geq a &\implies T_\alpha(\tilde{\rho}_0, \tilde{\rho}_1) \geq (1 - e^{-k})/2, \\ T_\alpha(\rho_0, \rho_1) \leq b &\implies T_\alpha(\tilde{\rho}_0, \tilde{\rho}_1) \leq 1/16. \end{aligned}$$

Here, the states  $\tilde{\rho}_0$  and  $\tilde{\rho}_1$  are defined over  $\tilde{O}\left(nk^{O\left(\frac{b \ln(2/a^2)}{a^2-b}\right)}\right)$  qubits. Moreover, when  $k \leq O(1)$  or  $a^2 - b \geq \Omega(1)$ , the time complexity of the procedure is  $\text{poly}\left(|Q_0|, |Q_1|, k, \exp\left(\frac{b \ln(1/a^2)}{a^2-b}\right)\right)$ .

In analogy with polarization lemmas for various classical [12, 21, 66] and quantum [51, 77] closeness measures, the proof of Lemma 25 proceeds by separately reducing the errors on both sides of the problem  $\text{QSD}_\alpha$ . This is achieved using the XOR lemma ([52, Lemma 5.4]) and the direct product lemma ([52, Lemma 5.5]) for  $T_\alpha$ . The full statements of these two lemmas, along with the full proofs of all three, are given in the full version, particularly in [52, Section 5.1].

## 5.2 Computational hardness and lower bounds for $\alpha > 1$ near 1

► **Theorem 26** ( $\text{QSD}_\alpha$  is QSZK-hard). *For any positive constant  $\delta > 0$  that can be made arbitrarily small, the following holds for sufficiently large  $n$ :*

- (1) *For any  $1 \leq \alpha \leq 1 + \frac{1}{n^{1+\delta}}$ ,  $\forall \tau \in (0, 1/2)$ ,  $\text{QSD}_\alpha[1 - \gamma_{\delta, \tau}(n), \gamma'_{\delta, \tau}(n)]$  is QSZK-hard, where  $\gamma_{\delta, \tau}(n) := 1 - 2^{-\frac{n+1}{n^{1+\delta}+1}} + 2^{-n^\tau - \frac{n+1}{n^{1+\delta}+1}}$  and  $\gamma'_{\delta, \tau}(n) := 2^{-n^\tau - \frac{1}{n^{1+\delta}+1}}$ .*
- (2) *For any  $1 + \frac{1}{n^{1+\delta}} < \alpha \leq 1 + \frac{1}{n}$ ,  $\forall \tau \in (0, 1/2)$ ,  $\text{QSD}_\alpha[\frac{1}{2} - 2^{-n^\tau-1}, 2^{-n^\tau - \frac{1}{n^{1+\delta}}}]$  is QSZK-hard.*

The proof of Theorem 26 (see also [52, Theorem 5.6]), based on reductions to  $\text{QSD}$  that is QSZK-hard [77, 78], is provided in the full version, particularly in [52, Section 5.2]. For any  $n$ -qubit state  $\rho$  of rank  $r$ , let  $\rho_{\text{U}}$  be the corresponding  $n$ -qubit state whose eigenvalues are uniformly distributed over the support of  $\rho$ . Next, we can establish the following lower bounds (see also [52, Theorem 5.7]), and the proof also can be found in [52, Section 5.2]:

► **Theorem 27** (Query complexity lower bounds for  $\text{QSD}_\alpha$ ). *The following query complexity lower bounds hold in the purified quantum query access model, depending on the range of  $\alpha$ , where  $\delta > 0$  is a constant that can be made arbitrarily small:*

- (1) *For any  $1 + \frac{1}{n^{1+\delta}} < \alpha \leq 1 + \frac{1}{n}$  and  $0 < \epsilon \leq 2^{\frac{1}{\alpha}-2}$ , there exist an  $n$ -qubit state  $\rho$  of rank  $r$  and the corresponding state  $\rho_{\text{U}}$  such that deciding whether  $T_\alpha(\rho, \rho_{\text{U}})$  is at least  $\epsilon$  or exactly 0 requires  $\Omega(r^{1/3})$  queries.*
- (2) *For any  $1 \leq \alpha \leq 1 + \frac{1}{n^{1+\delta}}$ , there exist a constant  $\epsilon > 0$  such that, for some  $n$ -qubit state  $\rho$  of rank  $r$  and the corresponding state  $\rho_{\text{U}}$ , estimating  $T_\alpha(\rho, \rho_{\text{U}})$  to within additive error  $\epsilon$  requires  $\tilde{\Omega}(r^{1/2})$  queries.*

By leveraging the same reduction used to prove Theorem 27(1), the rank-dependent sample complexity lower bound (given in [52, Lemma 2.10(2)] in the full version) for estimating the trace distance  $T(\cdot, \cdot)$  can be extended to the quantum  $\ell_\alpha$  distance  $T_\alpha(\cdot, \cdot)$  with  $1 \leq \alpha \leq 1 + \frac{1}{n}$ , as also stated in [52, Theorem 5.8]:

► **Theorem 28** (Sample complexity lower bound for  $\text{QSD}_\alpha$ ). *For any  $1 \leq \alpha \leq 1 + \frac{1}{n}$  and  $0 \leq \epsilon \leq 2^{\frac{1}{\alpha}-2}$ , there exists an  $n$ -qubit state  $\rho$  of rank  $r$  and the corresponding state  $\rho_{\text{U}}$  such that deciding whether  $T_\alpha(\rho, \rho_{\text{U}})$  is at least  $\epsilon$  or exactly 0 requires  $\Omega(r/\epsilon^2)$  samples of  $\rho$ .*

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